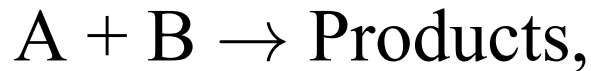


Collisional Processes

- Long range interaction --- between ions/electrons and ions/electrons; Coulomb $1/r$
- Intermediate range interaction --- between ions/electrons and neutral atoms/molecules; Induced dipole $1/r^4$
- Short range interaction --- between neutrals, $1/r^6$

In general, for a two-body collision,



the reaction rate per unit volume = $n_A n_B \langle \sigma v \rangle_{AB}$,

where the rate coefficient is

$$\langle \sigma v \rangle_{AB} = \int_0^{\infty} \sigma_{AB} v f(v) dv \quad [\text{cm}^3 \text{ s}^{-1}]$$

v = relative velocity between A and B

$\sigma_{AB}(v)$ = reaction cross section; velocity dependent

$f(v)$ = velocity distribution function

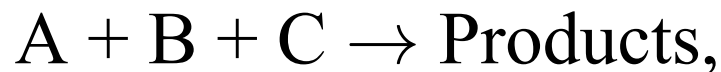
In thermal equilibrium

$$f_v dv = 4\pi \left(\frac{\mu}{2\pi kT} \right)^{3/2} e^{-\mu v^2 / 2kT} v^2 dv$$

In terms of energy

$$\langle \sigma v \rangle_{AB} = \left(\frac{8kT}{\pi\mu} \right)^{1/2} \int_0^\infty \sigma_{AB}(E) \frac{E}{kT} e^{-E/kT} \frac{dE}{kT}$$

If the density is high, e.g., in the Earth's atmosphere,
three-body collision may become important,



the reaction rate per unit volume = $k_{ABC} n_A n_B n_C$

where k_{ABC} is the three-body collisional rate
coefficient [$\text{cm}^6 \text{s}^{-1}$]

Elastic scattering by an inverse-square force, e.g., Rutherford scattering

Exact solutions complicated; use the “**impact approximation**”, i.e., motion in a straight line

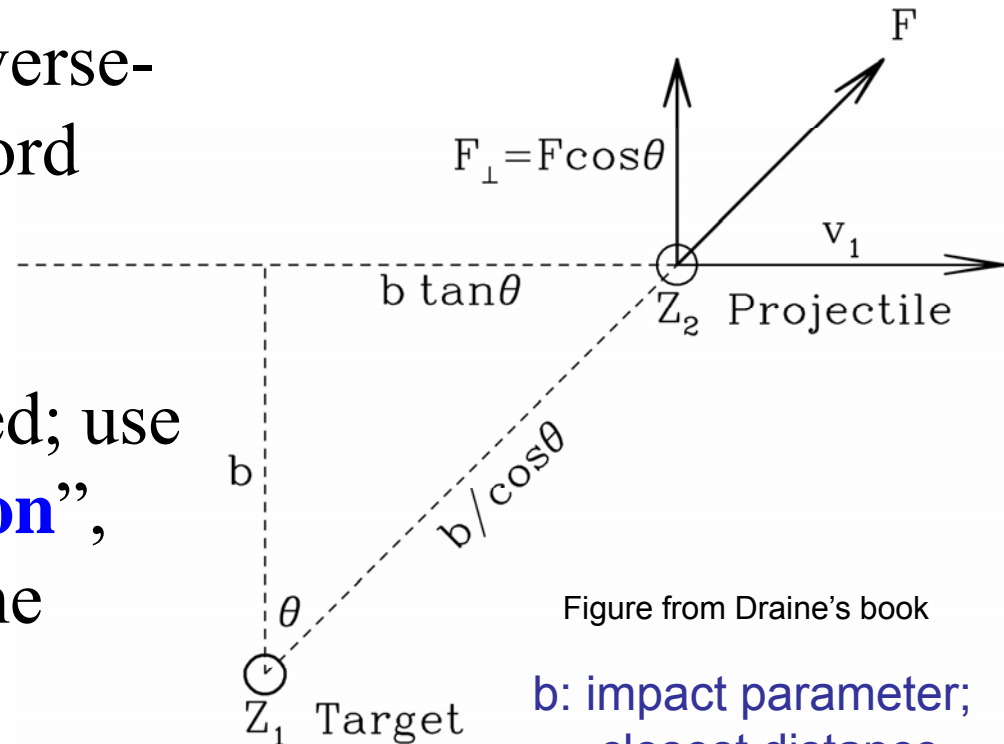


Figure from Draine's book

b : impact parameter;
closest distance
 v_1 : relative velocity

Assumption: constant velocity during the encounter between the target and the projectile

Question: How much momentum is transferred (\perp direction)?

Impact Approximation

Coulomb force

$$F_{\perp} = \frac{Z_1 e Z_2 e}{(b / \cos \theta)^2} \cos \theta = \frac{Z_1 Z_2 e^2}{b^2} \cos^3 \theta$$

$$\text{Interaction time scale } dt = \frac{d(b \tan \theta)}{v_1} = \frac{b}{v_1} \frac{d\theta}{\cos^2 \theta}$$

Total momentum transfer is

$$\begin{aligned} \Delta p_{\perp} &= \int_{-\infty}^{\infty} F_{\perp} dt = \frac{Z_1 Z_2 e^2}{b v_1} \int_{-\pi/2}^{\pi/2} \cos \theta d\theta \\ &= \frac{2 Z_1 Z_2 e^2}{b v_1} \approx \underbrace{\frac{Z_1 Z_2 e^2}{b^2}}_{\text{Force at closest distance}} \underbrace{\frac{b}{v_1}}_{\text{Time scale}} \end{aligned}$$

In a collisional ionization, there must be enough momentum transfer.
 $(E=p^2/2m)$

Fast moving
 $\rightarrow (1/2) m_e v^2 \gg E_I$
 $\oplus \quad e^-$

$$(\Delta P_{\perp})^2 > 2mE_I \Rightarrow \left(\frac{2Z_1Z_2e^2}{bv_1}\right)^2 > 2mE_I$$

So,

$$b^2 < b_{\max}^2(v) = \frac{(2Z_1Z_2e^2)^2}{v_1^2 \cdot 2mE_I} = \frac{2Z_p^2e^4}{m_ev^2E_I}$$

and the ionization cross section becomes

$$\sigma(v) \approx \pi b_{\max}^2 = \frac{2\pi Z_p^2e^4}{m_ev^2E_I} \quad \text{This is ok if } v \uparrow \uparrow.$$

For $v_{\min} = \frac{1}{2} m_e v_{\min}^2 = E_I$

$$\begin{aligned}
 \langle \sigma v \rangle &= \int \sigma(v) v f(v) dv \\
 &= \int_{v_{\min}}^{\infty} \frac{2\pi Z_p^2 e^4}{m_e v^2 E_I} v 4\pi \left(\frac{m_e}{2\pi kT}\right)^{3/2} v^2 e^{-m_e v^2 / 2kT} dv \\
 &= Z_p^2 \left(\frac{8\pi}{m_e kT}\right)^{1/2} \frac{e^4}{E_I} e^{-E_I/kT}
 \end{aligned}$$

For an H atom at level n , $E_I = 13.6 \text{ [eV]}/n^2$, so for a large n , e.g., $n \sim 100$, and $T \sim 10^4 \text{ K}$, $E_I \downarrow \downarrow (\ll kT)$
 \rightarrow in radio frequencies.

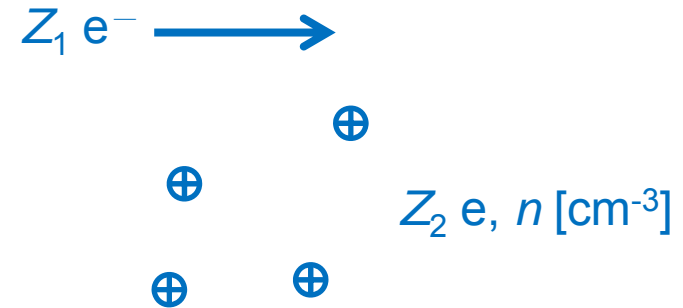
$\langle \sigma v \rangle \propto \frac{1}{E_I} \propto n^2$, so is very large.

Deflection Timescale

Net momentum transfer

$$\left\langle \frac{d}{dt} (\Delta P_{\perp})^2 \right\rangle$$

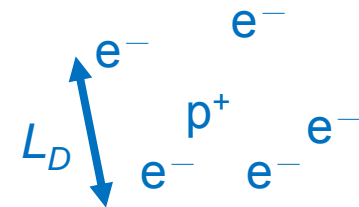
There must be a range of distance, for which
 $b_{\min} = Z_1 Z_2 e^2 / \text{Energy}$, and
 $b_{\max} \approx L_D$ (Debye length)



In plasma, the distributions of ions and electrons are correlated because of charge neutrality.

Near a proton \rightarrow more electrons than protons
 \rightarrow the proton is “shielded”

Average charge within a region $\langle Q(L_D) \rangle = -e$



$$L_D = \left(\frac{kT}{4\pi n_e e^2} \right)^{1/2} = 690 T_4^{1/2} \left[\frac{n_e}{\text{cm}^{-3}} \right]^{-1/2} [\text{cm}]$$

$$\langle \frac{d}{dt}(\Delta P_{\perp})^2 \rangle \propto \frac{n_2}{v_1} \ln \Lambda$$

So Λ is large, $\equiv b_{\max}/b_{\min} \approx 20 - 35$ in ISM conditions.

\Rightarrow For elastic scattering of electrons by ions, weak distant encounters (\gg atomic scales) more important than close encounters.

If an electron comes in in \approx atomic dimensions,
the atom is

suddenly perturbed \rightarrow transition

\rightarrow deexcitation \rightarrow line radiation

$$\langle \sigma v \rangle_{10} = \frac{8.629 \times 10^{-8}}{\sqrt{T_4}} \frac{\Omega_{10}}{g_1} [\text{cm}^3 \text{s}^{-1}]$$

where Ω_{10} is **collision strength**.

(i) almost independent of T for $T < T_4 = 10^4$ K

(ii) $1 < \Omega_{10} < 10$

Ion-Neutral Collisions

Neutral is polarized.



Dipole moment $\vec{P} = \alpha_N \vec{E}$

Interaction potential

$$U(r) = -\frac{1}{2}\alpha_N \frac{Z^2 e^2}{r^4} \propto r^{-4}$$

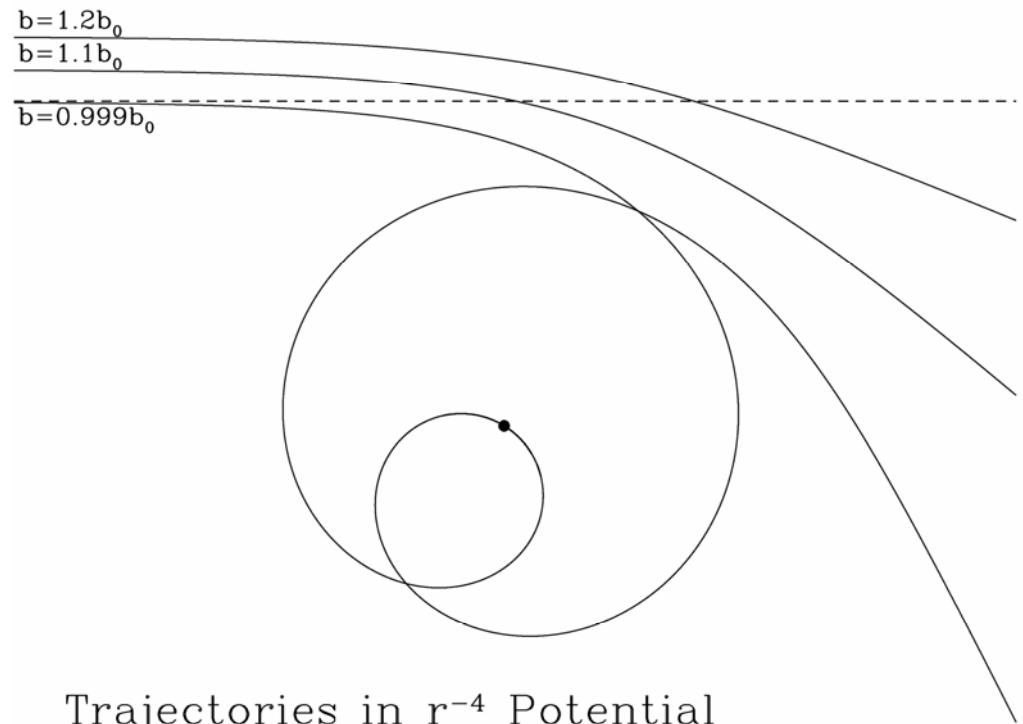
where α_N is **polarizability** \approx a few a_0 .

$$a_0 = \text{Bohr radius} \equiv \frac{\hbar^2}{m_e e^2} = 5.292 \times 10^{-9} \text{ [cm]}$$

For such a potential, if $b < b_0$, the deflection cross section is large.

$\sigma = \pi b_0^2 \propto 1/v$, and the rate coefficient $\langle \sigma v \rangle \propto T$

Usually if $T \uparrow$, $\sigma \downarrow$



The ion-neutral reactions are important in cool ISM.

Electron-Neutral Collisions

In low-ionization ISM (e.g., protoplanetary disks) ions are rare. e^- -neutral (H_2 , He) scattering is important.

e^- - H_2 scattering

- (1) If $E < 0.044$ eV \rightarrow pure elastic scattering
- (2) If $E > 0.044$ eV \rightarrow rotational excitation possible
- (3) If $E > 0.5$ eV \rightarrow vibrational excitation possible
- (4) If $E > 11$ eV \rightarrow electronic excitation

By experiment $\sigma \simeq 7.3 \times 10^{-6} \left(\frac{E}{0.01 \text{ eV}} \right) [\text{cm}^2]$

and $\langle \sigma v \rangle \simeq 4.8 \times 10^{-9} \left(\frac{T}{10^2 \text{ K}} \right)^{0.68} [\text{cm}^3 \text{ s}^{-1}]$

Neutral-Neutral Collisions

Repulsive if distance ↓

Weakly attractive if distance ↑ ∴ van der Waals interactions (mutual induced electric dipole)

$$U(r) \propto r^{-6}$$

Hard sphere OK; radii $R \sim 1\text{\AA}$

$$b < R_1 + R_2 \quad \sigma = \pi (R_1 + R_2)^2 \sim 1.2 \times 10^{-15} [\text{cm}^2]$$

$$\langle \sigma v \rangle = 1.81 \times 10^{-10} \left(\frac{T}{10^2 \text{K}} \right)^{1/2} \left(\frac{m_H}{\mu} \right)^{1/2} \left(\frac{R_1 + R_2}{2\text{\AA}} \right)^2 [\text{cm}^3 \text{s}^{-1}]$$

Collision

Gas (hydrogen atoms) root-mean-squared speed

$$m_{\text{H}} \sqrt{\langle v^2 \rangle} = 3kT$$

For H I regions, $T \sim 100 \text{ K}$, $\langle v \rangle_{\text{HI}} \sim 1 \text{ km s}^{-1}$

For e^- , $\langle v \rangle_{e^-} \sim 50 \text{ km s}^{-1}$

Cross sections σ

- Hard sphere OK for neutral atoms, $\sigma = \pi(a_1 + a_2)^2$
i.e., ‘physical’ cross section

$$\sigma_{\text{HI,HI}} \leftarrow a \sim 5.6 \times 10^{-9} \text{ cm}$$

$$\text{c.f., Bohr radius (first orbit)} = 5.3 \times 10^{-9} \text{ cm}$$



Cross sections σ

- For free e^- , p^+

$\sigma \gg \sigma_{\text{physical}}$ because of Coulomb force, need QM

$$a \sim \frac{2.5 \times 10^{-2}}{v^2} \text{ cm } (v \text{ in km})$$

If $v_{e^-} \sim 50 \text{ km s}^{-1}$, $a \sim 10^{-5} \text{ cm}$ for e^- - e^- collision

$$T = 3 \times 10^4 \text{ K}, \langle v \rangle \sim 10^3 \text{ km s}^{-1}$$

$$\longrightarrow a \sim 2.5 \times 10^{-8} \text{ cm}$$

c.f., classical electron radius $\sim 2.8 \times 10^{-13} \text{ cm}$

Conventional unit for cross section

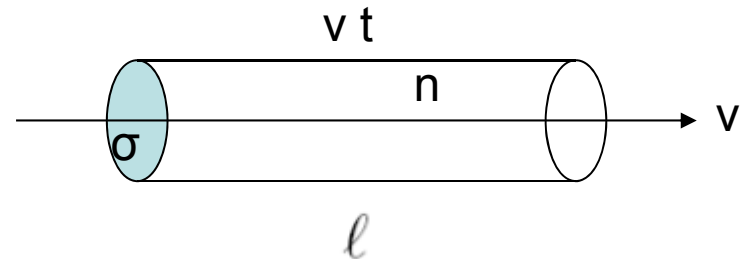
$$1 \text{ barn} = 10^{-24} \text{ cm}^2$$

$$\sigma_{\text{HI,HI}} \sim 10^8 \text{ barns } (\sim 10^{-16} \text{ cm}^2)$$

$$\frac{e^2}{r_0} = m c^2$$

$$r_0 = \frac{e^2}{m c^2} \sim 2.8 \times 10^{-13} \text{ cm}$$

Collision



of collisions = # of particles in the (moving) volume

$$N = n \sigma v t$$

of collisions per unit time = $N/t = n \sigma v$

Time (mean-free time) between 2 consecutive collisions ($N=1$) = $t_{\text{collision}} = \frac{1}{n \sigma v}$

Mean-free path $\ell = v t_{\text{collision}}$, i.e., $\ell = \frac{1}{n \sigma}$

Ex 1 $n_{\text{HI}} \sim 10 \text{ cm}^{-3}; v_{\text{HI}} \sim 1 \text{ km s}^{-1}; \sigma_{\text{HI,HI}} \sim 10^{-16} \text{ cm}^2$

$$t_{\text{HI,HI}} \sim 10^{10} \text{ s} \sim 300 \text{ years}$$

$$\ell \sim 10^{15} \text{ cm} \sim 100 \text{ AU}$$

\therefore Collisions are indeed very rare.

Ex 2 $\sigma_{\text{e}^-, \text{HI}} \sim 10^{-15} \text{ cm}^2$ (polarization)

$$t_{\text{e}^-, \text{HI}} \sim \frac{1}{10 \times 10^{-15} \times 10^5} \sim 30 \text{ years}$$

Ex 3 $\sigma_{\text{e}^-, \text{e}^-} \sim 10^{-12} \text{ cm}^2; n_e \sim 0.2 \text{ cm}^{-3}$

$$t_{\text{e}^-, \text{e}^-} \sim \frac{1}{0.2 \times 10^{-12} \times 50 \times 10^5} \sim 10 \text{ days}$$