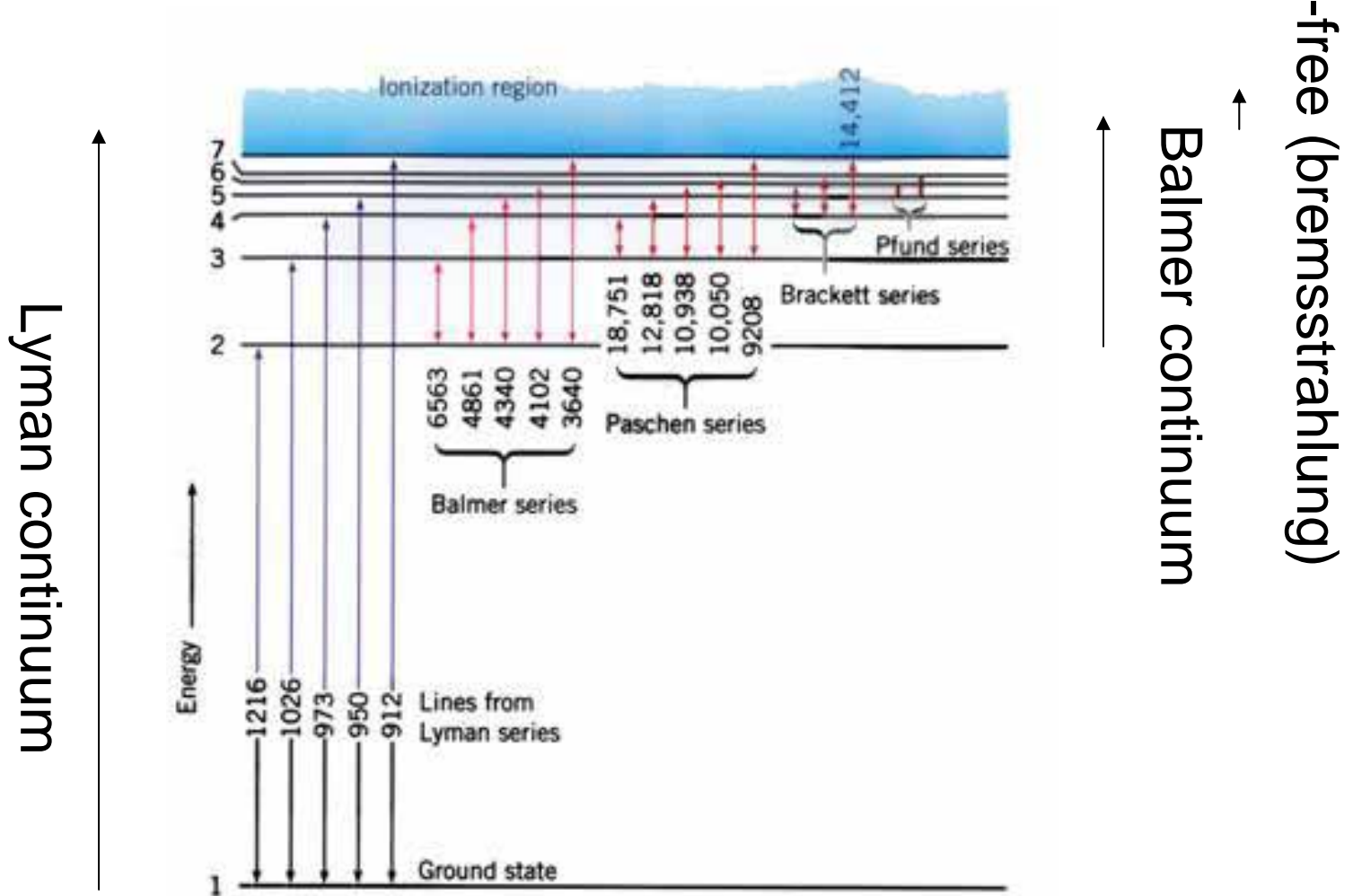
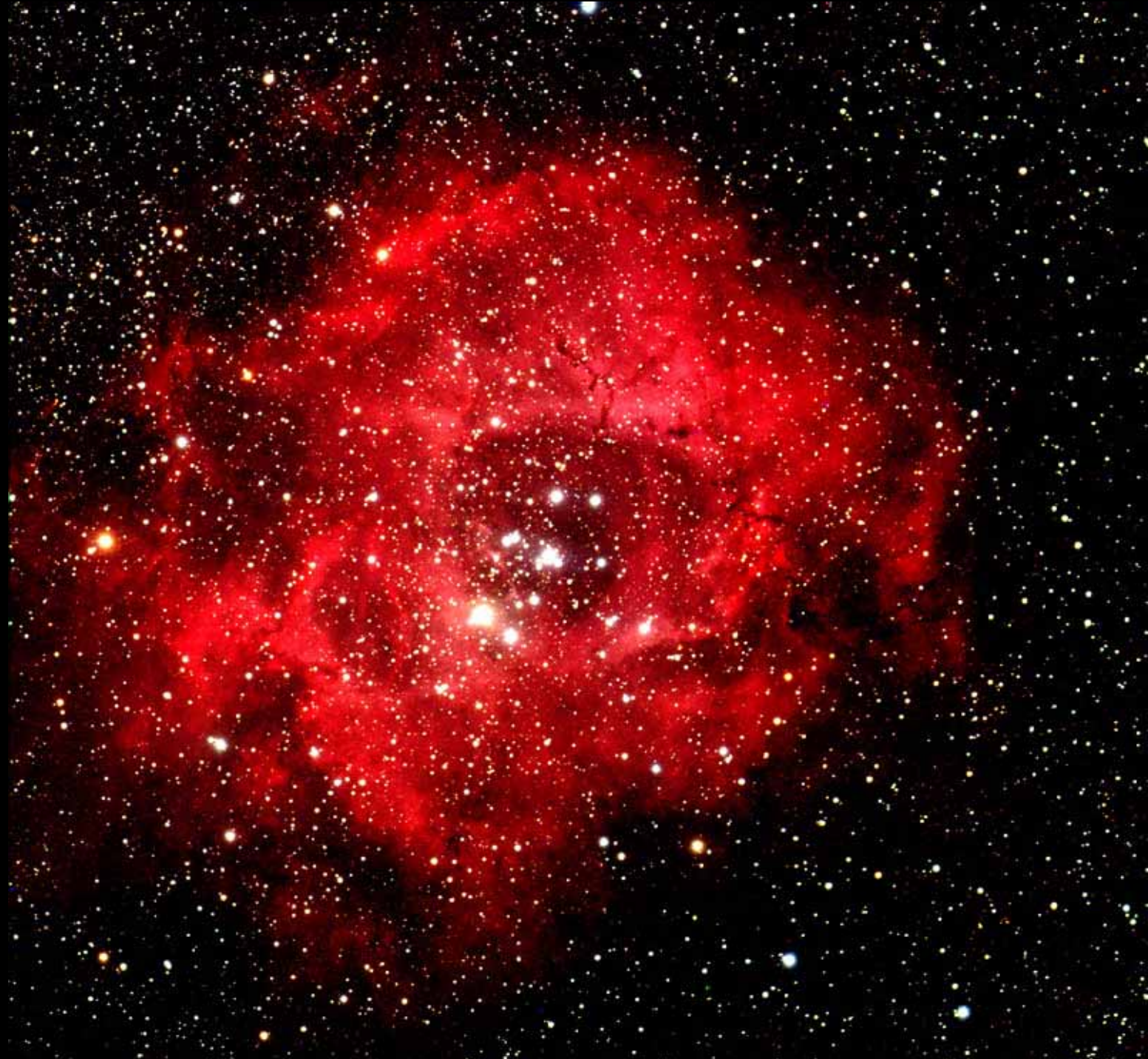
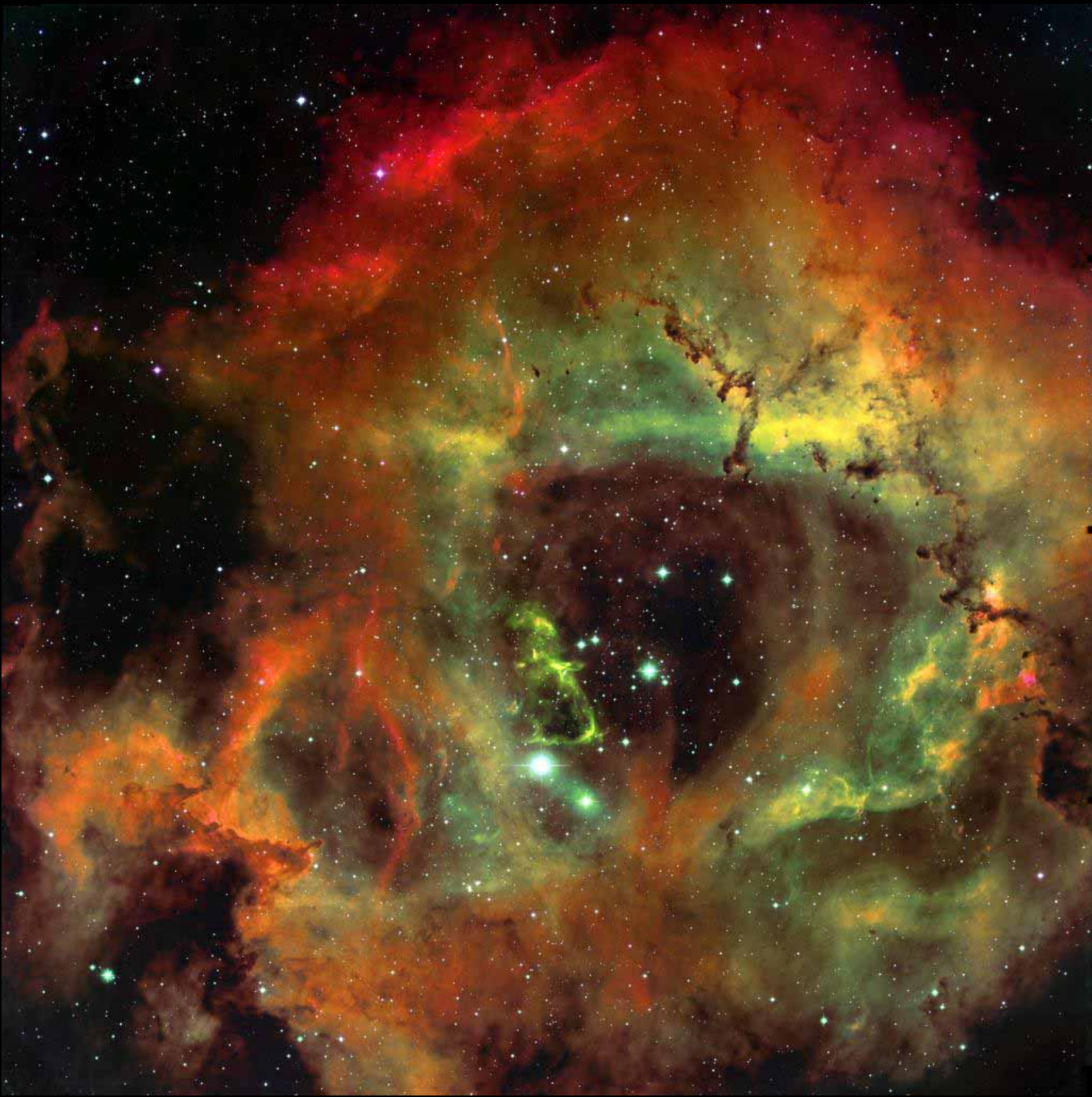


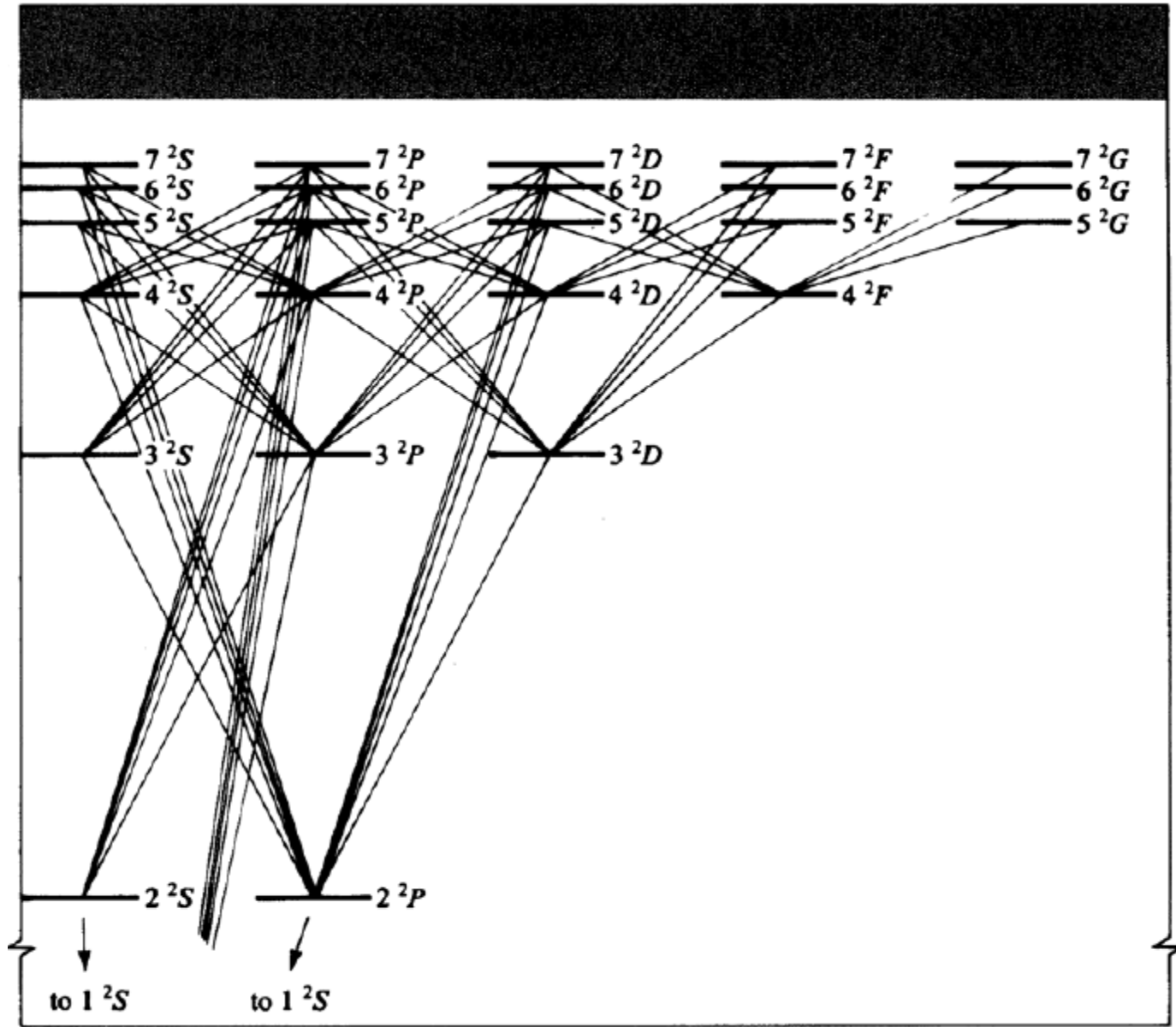
# Photoionization







# Hydrogen spectrum --- permitted transitions

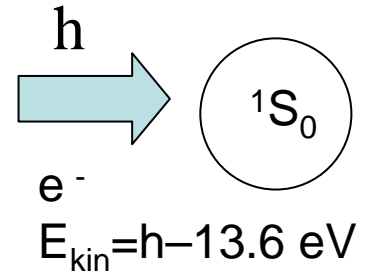


# Wavelengths of important H Lines

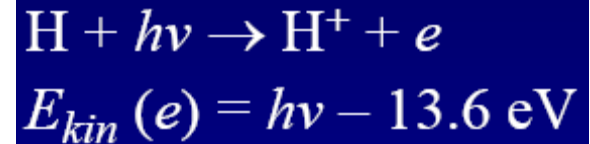
- Ly $\alpha$ :  $\lambda_{vac}$  = 1215.68 Å (space UV)
  - H $\alpha$ :  $\lambda_{air}$  = 6562.73 Å (red) 3-2
  - H $\beta$ :  $\lambda_{air}$  = 4861.33 Å (blue) 4-2
  - H $\gamma$ :  $\lambda_{air}$  = 4340.47 Å (blue) 5-2
  - H $\delta$ :  $\lambda_{air}$  = 4101.47 Å (violet) 6-2
  - Pa $\alpha$ :  $\lambda_{air}$  = 1.875  $\mu$ m (poor transmission)
  - Br $\alpha$ :  $\lambda_{air}$  = 4.051  $\mu$ m (difficult)
  - Br $\gamma$ :  $\lambda_{air}$  = 2.166  $\mu$ m (in infrared K band)
- } Balmer lines

- Transition probabilities between upper state  $u$  and lower state  $l$  are characterized by Einstein  $A$  and  $B$  coefficients or related oscillator strength  $f$  (see appendix for def)
  - Absorption  $l \longrightarrow u$ : oscillator strength  $f_{lu}$
  - Emission  $u \longrightarrow l$ : spontaneous emission coeff:  $A_{ul}$ 
    - For transition in H:  $A_{ul} \approx 10^4 \dots 10^8 \text{ s}^{-1}$

$$E_{\text{H,ion}} = 13.6 \text{ eV} \quad (\lambda = 912\text{\AA})$$



Probability of **photoionization** →  
photoionization cross section



For hydrogen-like atoms, the cross section is

$$\sigma_{\nu}^{\text{ion}} = \frac{7.9 \times 10^{18}}{Z^2} \left(\frac{\nu_1}{\nu}\right)^3 g_{1f} [\text{cm}^2], \quad \text{for } \nu > \nu_1$$

where  $g_{1f}$  is Gaunt factor  $\approx 1$  at optical wavelengths,  
 $h\nu_1 = Z^2 h\nu_0 = 13.6 Z^2 \text{ eV}$

For hydrogen,  $\nu_1 = 3.29 \times 10^{15}$  Hz,  $g_{1f} \approx 0.8$ , and a good approximation,

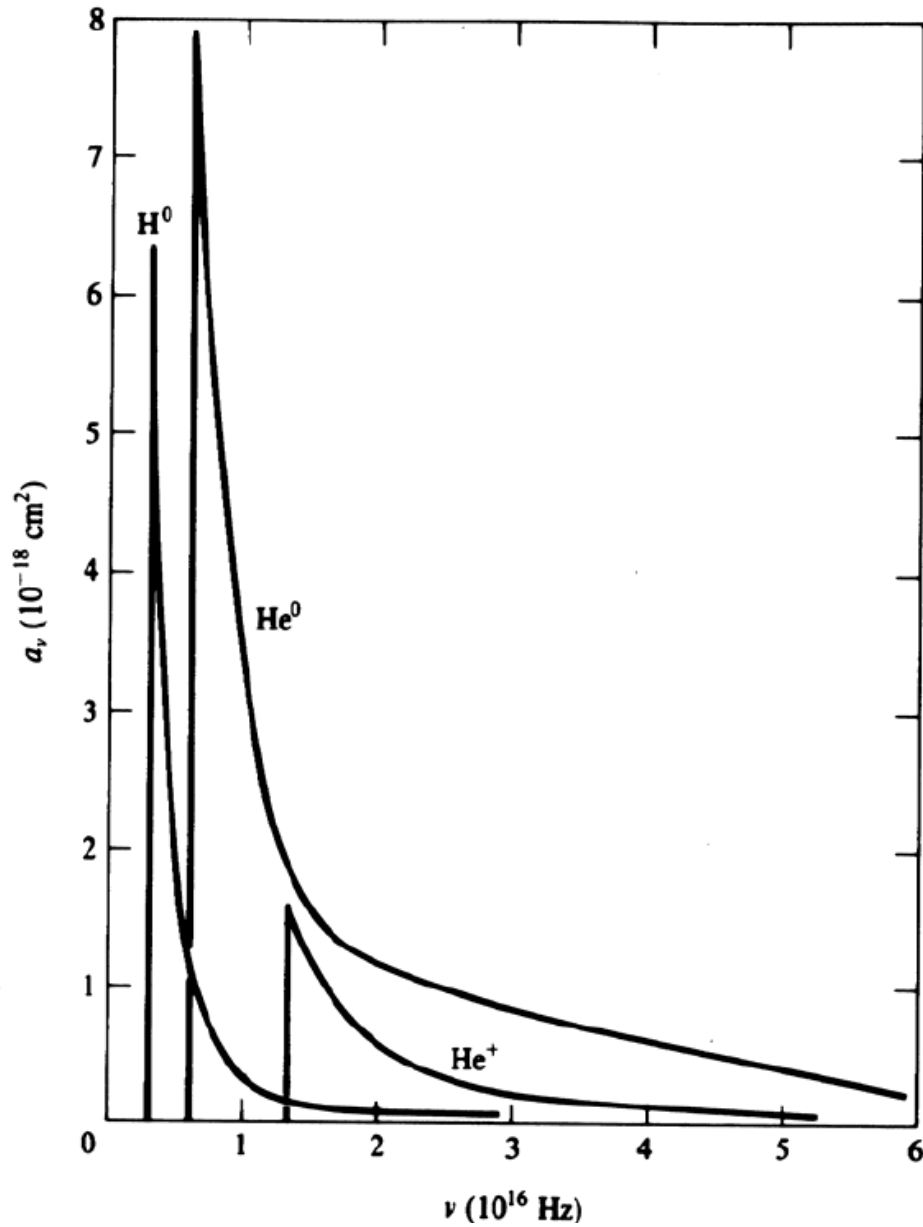
$$\sigma_{\text{ion}}(\nu) \approx 6.3 \times 10^{-18} \left(\frac{\nu_1}{\nu}\right)^3 [\text{cm}^2]$$

That is, high-energy photons, with much smaller photoionization absorption cross sections, penetrate deeper into neutral gas before being absorbed.

$$\begin{aligned} & [\# \text{ of ionization}] \text{ s}^{-1} \text{ atom}^{-1} \text{ due to photons in } \nu \text{ to } \nu + d\nu \\ &= \sigma_\nu \frac{4\pi \bar{I}_\nu d\nu}{h\nu} \end{aligned}$$



# Photoionization Cross Sections for H, He and He<sup>+</sup>



Define the coefficient  $\alpha$ , so that

$$\alpha n_e n_p = [\# \text{ of recombinations}] \text{ s}^{-1} \cdot \text{cm}^{-3}$$

$$\alpha = \langle v \sigma_{\text{recomb}} \rangle \quad \alpha(n, L) = \int v \sigma_{nK} f(\vec{v}) d^3 \vec{v}$$

But recombination may end up at different levels

$$\alpha^{(n)} = \sum_{m=n}^{\infty} \alpha_m$$

$\alpha^{(1)}$ : total recombination coefficient summed over all levels

$\alpha^{(2)}$ : total recombination coefficient excluding captures to  $n = 1$  level

$\alpha$ s can be computed exactly for hydrogen.

$$\alpha^{(1)} = \sum_{n=1}^{\infty} \alpha_n = 6.82 \times 10^{-13} \text{ cm}^3 \cdot \text{s}^{-1} \text{ (at 5000 K)}$$

$$\alpha^{(2)} = \sum_{n=2}^{\infty} \alpha_n = 4.54 \times 10^{-13} \text{ cm}^3 \cdot \text{s}^{-1} \text{ (at 5000 K)}$$

Spitzer gives  $\alpha^{(2)} = 2.59 \times 10^{-3} T_4^{-0.81}$

TABLE 2.1  
*Recombination coefficients<sup>a</sup>  $\alpha_n \ ^2L$  for H*

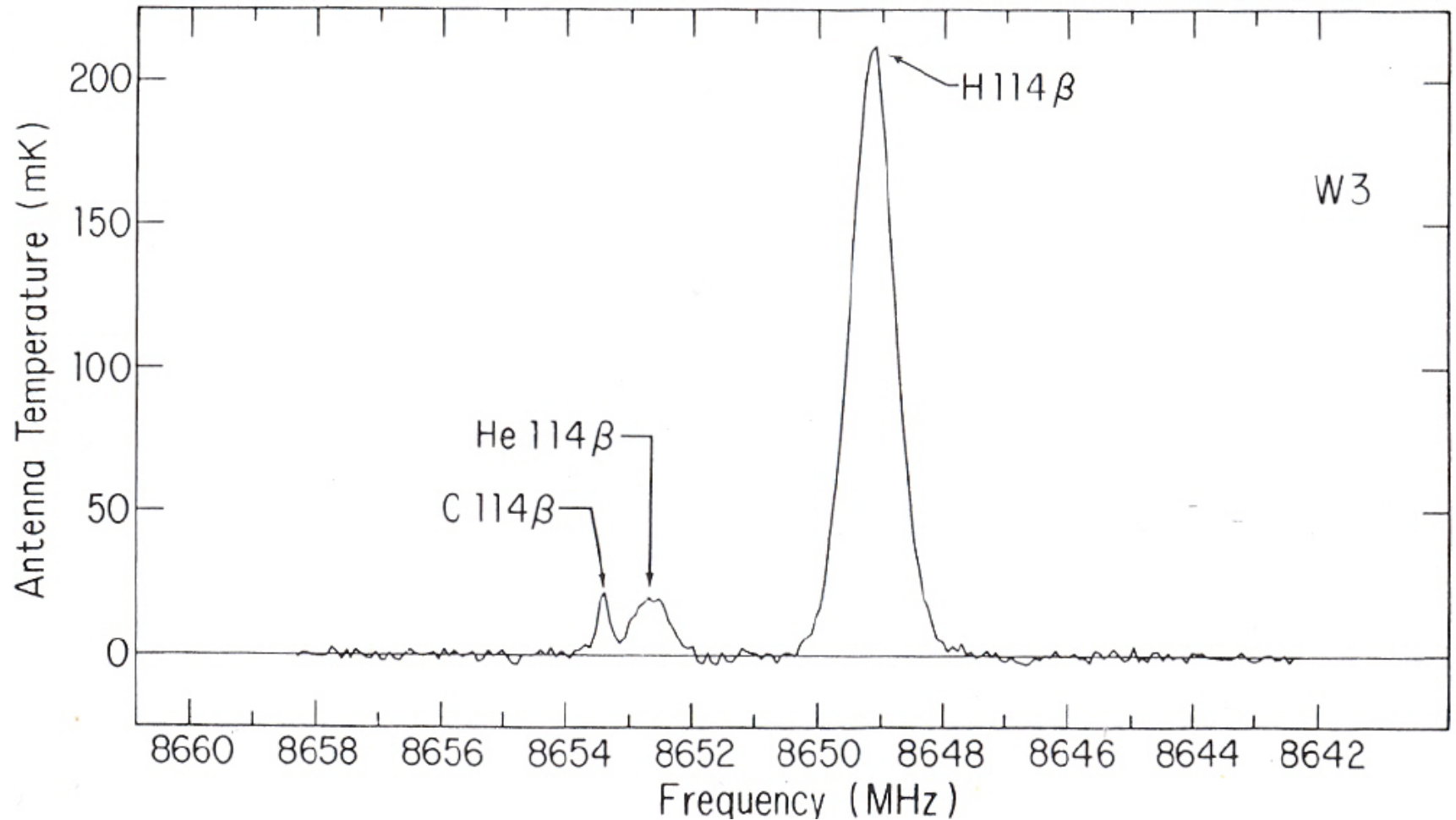
$\alpha_n \ ^2L$	<i>T</i>		
	5000° K	10,000° K	20,000° K
$\alpha_1 \ ^2S$	$2.28 \times 10^{-13}$	$1.58 \times 10^{-13}$	$1.08 \times 10^{-13}$
$\alpha_2 \ ^2S$	$3.37 \times 10^{-14}$	$2.34 \times 10^{-14}$	$1.60 \times 10^{-14}$
$\alpha_2 \ ^2P$	$8.33 \times 10^{-14}$	$5.35 \times 10^{-14}$	$3.24 \times 10^{-14}$
$\alpha_3 \ ^2S$	$1.13 \times 10^{-14}$	$7.81 \times 10^{-15}$	$5.29 \times 10^{-15}$
$\alpha_3 \ ^2P$	$3.17 \times 10^{-14}$	$2.04 \times 10^{-14}$	$1.23 \times 10^{-14}$
$\alpha_3 \ ^2D$	$3.03 \times 10^{-14}$	$1.73 \times 10^{-14}$	$9.09 \times 10^{-15}$
$\alpha_4 \ ^2S$	$5.23 \times 10^{-15}$	$3.59 \times 10^{-15}$	$2.40 \times 10^{-15}$
$\alpha_4 \ ^2P$	$1.51 \times 10^{-14}$	$9.66 \times 10^{-15}$	$5.81 \times 10^{-15}$
$\alpha_4 \ ^2D$	$1.90 \times 10^{-14}$	$1.08 \times 10^{-14}$	$5.68 \times 10^{-15}$
$\alpha_4 \ ^2F$	$1.09 \times 10^{-14}$	$5.54 \times 10^{-15}$	$2.56 \times 10^{-15}$
$\alpha_{10} \ ^2S$	$4.33 \times 10^{-16}$	$2.84 \times 10^{-16}$	$1.80 \times 10^{-16}$
$\alpha_{10} \ ^2G$	$2.02 \times 10^{-15}$	$9.28 \times 10^{-16}$	$3.91 \times 10^{-16}$
$\alpha_{10} \ ^2M$	$2.7 \times 10^{-17}$	$1.0 \times 10^{-17}$	$4. \times 10^{-18}$
$\alpha_A$	$6.82 \times 10^{-13}$	$4.18 \times 10^{-13}$	$2.51 \times 10^{-13}$
$\alpha_B$	$4.54 \times 10^{-13}$	$2.59 \times 10^{-13}$	$2.52 \times 10^{-13}$

<sup>a</sup> In  $\text{cm}^3 \text{sec}^{-1}$ .

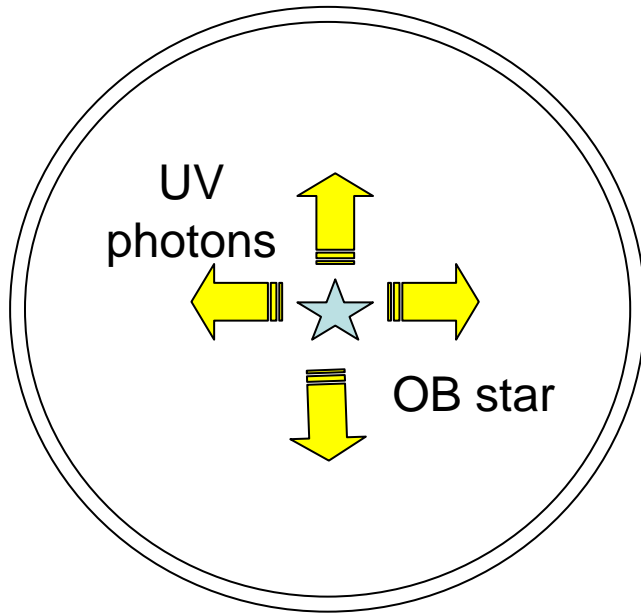
Recombination  $\rightarrow$  photons

Some of the photons can ionize or excite other species

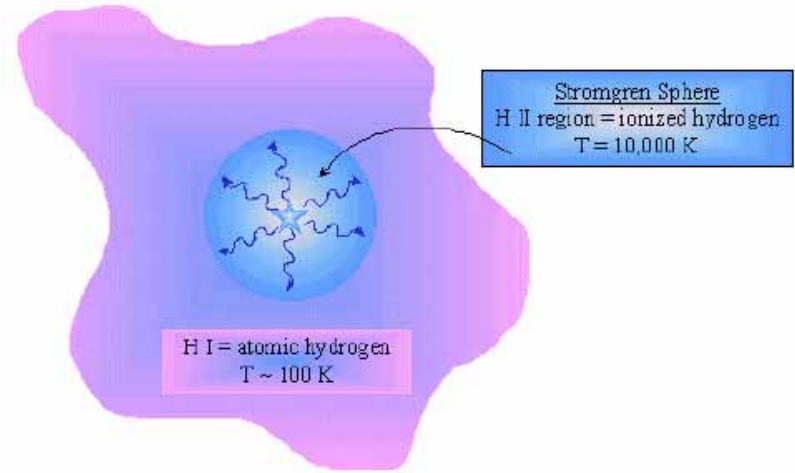
H 114 :  $n=116 \rightarrow 114$



# Application: H II regions



Once ionized, e recombines with p, emitting Balmer, Paschen, Pfund lines and continua



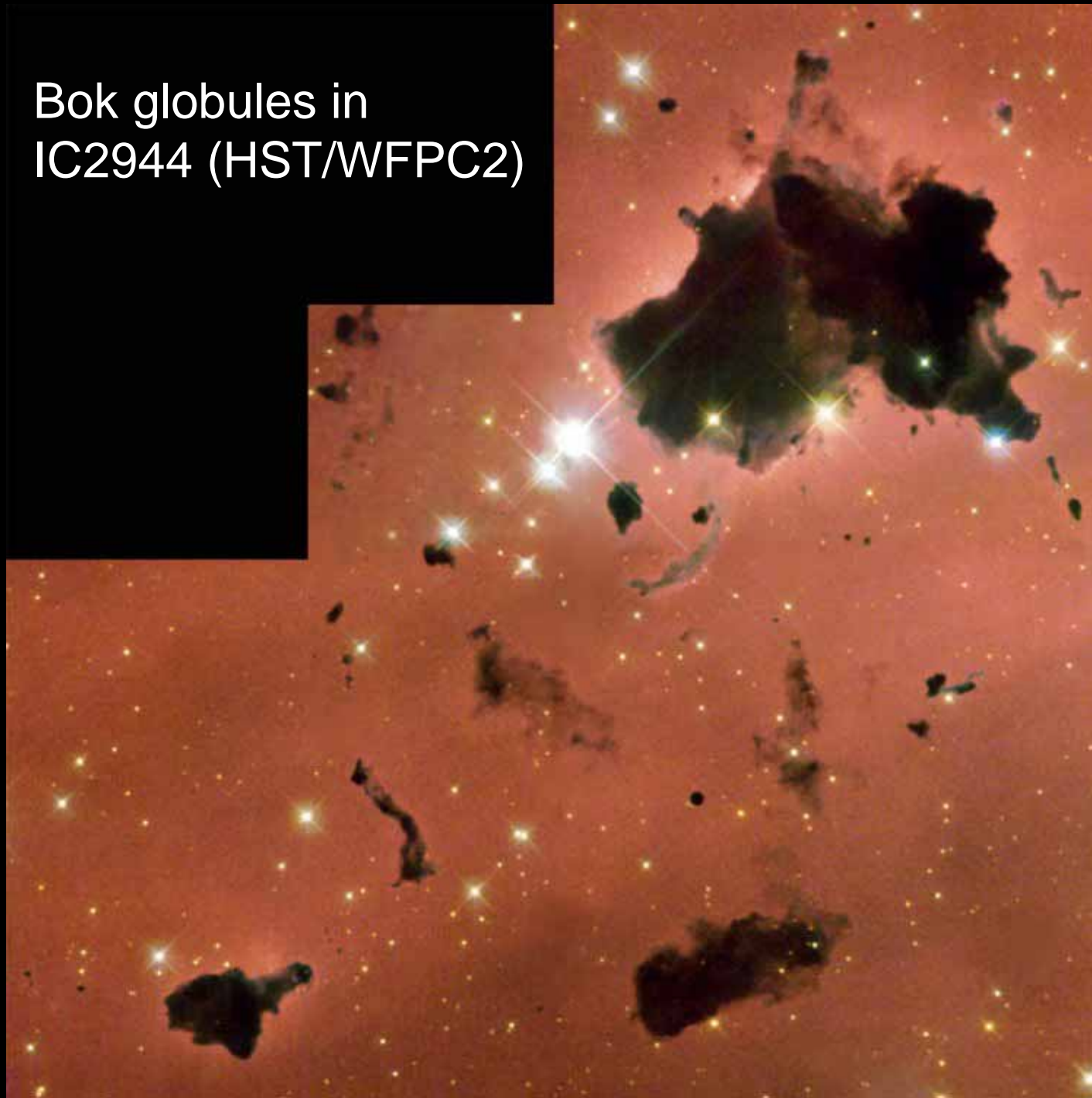
Radiation  $< 912\text{\AA}$   $\rightarrow$  ionization from gr state

If e already in an excited state, a longer will do

Collisional ionization negligible in HII regions

$\longrightarrow$  e cascading  $\rightarrow$  H

Bok globules in  
IC2944 (HST/WFPC2)



**Strings of red H II regions delineate the arms of the Whirlpool Galaxy.**



# Strömgren Sphere

The Strömgren radius,  $R_s$ , within which  
total # of recombinations to levels except the gr. state  
= total # of ionizing photons from the luminous star

Total recombinations:  $\alpha^{(2)} n_e n_p (4\pi R_s^3 / 3)$

Total stellar ionizing photons  $\nu > \nu_0$ :  $\int_{\nu_0}^{\infty} (L_\nu / h\nu) d\nu$

$L_\nu$ : stellar luminosity at  $\nu$  [ergs s<sup>-1</sup> Hz<sup>-1</sup>]

$$4\pi R_*^2 \int_{\lambda=912}^{\infty} \frac{\pi B_\nu(R_*)}{h\nu} d\nu = \alpha^{(2)} n_e n_p (4\pi R_s^3 / 3)$$



Within  $R_s$ , ionization is complete,  $n_e \approx n_p \approx n_H$

Outside  $R_s$ ,  $n_e \approx n_p \approx 0$

Designate the LHS (# of Lyman photons) as  $L_{912}^*$ , we get

$$L_{912}^* = 4\pi R_s^3 \alpha^{(2)} n_H^2 / 3$$

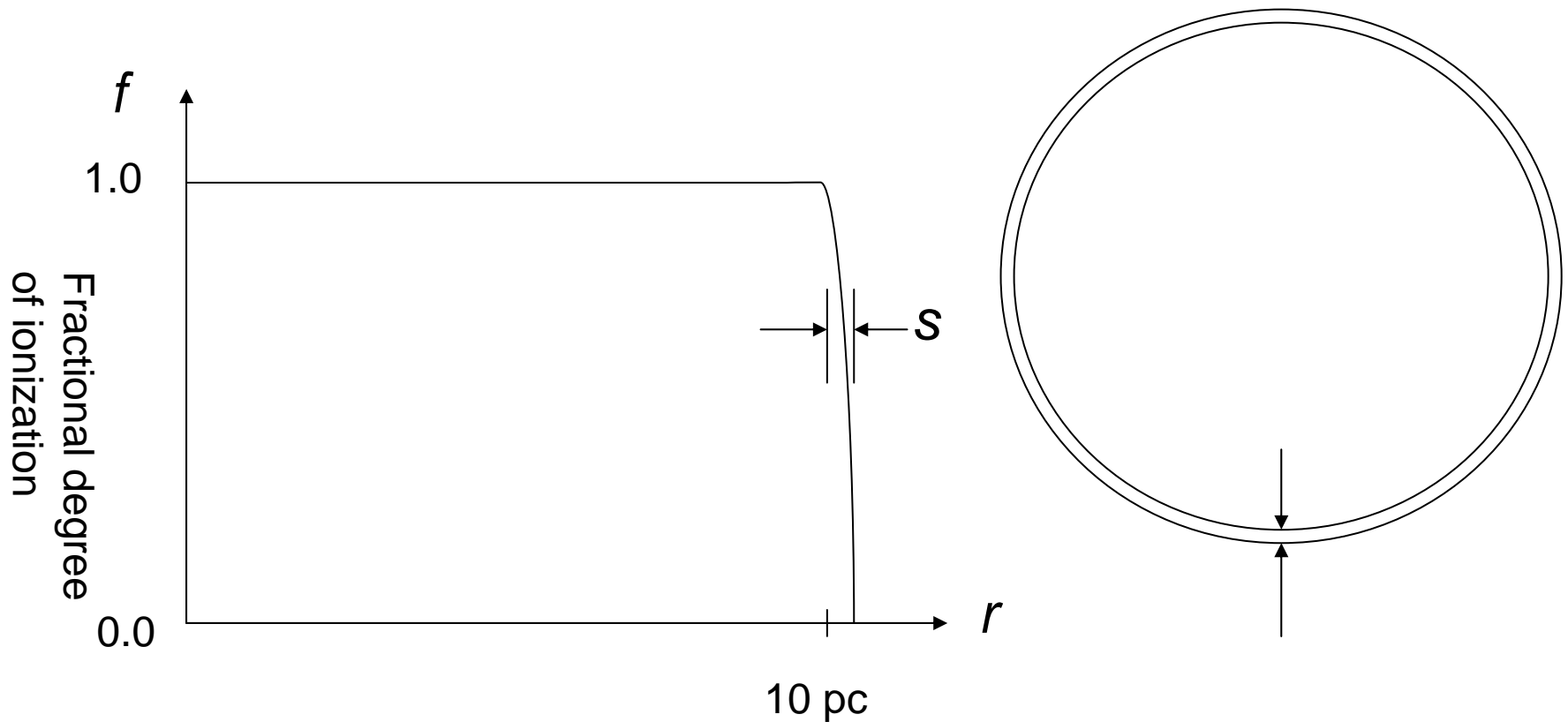
$$R_s = \left( \frac{3L_{912}^*}{4\pi\alpha^{(2)}n_H^2} \right)^{1/3}$$

$$L_{912}^* \approx 5 \times 10^{42} \exp(3.2 \times 10^{-4} T_*)$$

$$R_s \approx 0.62 n_H^{-2/3} \exp(1.07 \times 10^{-4} T_*) \text{ pc}$$

Take  $T_* = 40,000$  K (i.e., O6 V),  $n_H = 10$ ,  $R_s \approx 1$  pc,  
then  $R_s \approx 10$  pc

How thick is the transition zone ( $s$ )?



In the transition zone,  $\tau = \sigma_{\nu}^{\text{ion}} n_H s \approx 1$

$$s = \frac{1}{\sigma_{\nu} n_H}$$

Given  $\sigma_{\nu_{912}} = 6.3 \times 10^{-18} \text{ cm}^2$ ,  $n_H = 10 \text{ cm}^{-3}$ ,  
 $s \approx 0.005 \text{ pc} \ll 10 \text{ pc}$

The boundary of an H II region is very sharp!

Spectral Type	$T_*$ [K]	$R_s (n_e n_p)^{1/3}$ [pc cm <sup>-2</sup> ]
O5	47,000	110
O9	34,500	38
B1	22,600	4.4

Note: The above assumes no dust absorption;  
otherwise  $R_s$

$A_v$ (mag)	0.1	0.5	1.0	2.0	5.0	10.0
$R'_s/R_s$	0.91	0.70	0.56	0.42	0.25	0.15

- He can also be ionized,  $\lambda < 506$  and  $\lambda < 208 \text{ \AA}$
- Stars very hot,  $T^* > 10^5 \text{ K}$
- Ionization Application --- **Zanstra Method**  
 If all  $\lambda < 912 \text{ \AA}$  photons absorbed by H atoms in the nebula  
 Each UV photon  
 → 1 Ly photon (absorbed and re-emitted)  
 + 1 Balmer photon (escaped readily)
- So, # of Balmer photons = # of  $\lambda < 912 \text{ \AA}$  photons → all energy radiated by the star  
 →  $T_*$

